

Training.

Relaxation finite element schemes for the incompressible Navier-Stokes equations

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16. Juli 09

Outline

- 1 Introduction
 - Navier-Stokes
 - Relaxation Finite Element Schemes for The Incompressible Navier-Stokes equations.
 - Implementation and Software.
 - Methodology
- 2 Stokes Equations
- 3 Steady State Navier-Stokes Equations
- 4 Incompressible Navier-Stokes Equations
- 5 Navier-Stokes equations with relaxation

Incompressible Navier-Stokes Equations

$$\begin{cases} \partial_t u + u \cdot \nabla u - \nu \Delta u + \nabla p & = f(x, t) \text{ on } \Omega \times [0, T] \\ \operatorname{div} u & = 0 \\ u|_{\partial\Omega} & = g(x, t) \text{ on } \partial\Omega \times [0, T] \end{cases} \quad (1)$$

- $\mathbb{R}^2 \supset \Omega$ space domain
- $[0, T] \subset \mathbb{R}$ time domain
- $u : \Omega \times [0, T] \longrightarrow \mathbb{R}^2, (x, t) \mapsto u(x, t)$ velocity
- $p : \Omega \times [0, T] \longrightarrow \mathbb{R}, (x, t) \mapsto p(x, t)$ pressure
- $\operatorname{div} u = 0$ describes incompressibility of the fluid.
- Given: Ω, T, f, g, ν
- Find: velocity u and pressure p

Navier-Stokes Equations with Relaxation.

$$\begin{cases} \partial_t u + u \cdot \nabla u - \nu \Delta u + \nabla p + \varepsilon \partial_{tt} u + \varepsilon \partial_t \nabla p & = f(x, t) \\ \operatorname{div} u & = 0 \\ u|_{\partial\Omega} & = g(x, t) \end{cases} \quad (2)$$

- $u : \Omega \times [0, T] \longrightarrow \mathbb{R}^2, (x, t) \mapsto u(x, t)$ velocity
- $p : \Omega \times [0, T] \longrightarrow \mathbb{R}, (x, t) \mapsto p(x, t)$ pressure
- Given: Ω, T, f, g, ν
- Find: u and p .

Advantage

Better description of the physics

Implementation and Software

Open Source and Free Software:

- Numerics: FreeFem++.
- Visualization: FreeFem++, gnuplot, R, (OpenDX)
- IDE: *FreeFem++-cs*, *Eclipse*.
- Publishing: \LaTeX , LyX, latex-beamer, Inkscape, TextText.
- Version Control System: subversion.
- Programming Languages: FreeFem++, python, C++, R.
- Programming paradigm: Literate Programming

Methodology

- Fix the goal. Get Familiar with the problem.
- Collect Information. Find the starting point.
- Go step by step increasing the complexity of the problem.

$$\begin{cases} \partial_t u + u \cdot \nabla u - \nu \Delta u + \nabla p + \varepsilon \partial_{tt} u + \varepsilon \partial_t \nabla p & = f(x, t) \\ \operatorname{div} u & = 0 \\ u|_{\partial\Omega} & = g(x, t) \end{cases}$$

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Stokes equations

$$\begin{cases} -\nu \Delta \tilde{u} + \nabla p & = \tilde{f}(x, t) \\ \operatorname{div} \tilde{u} & = 0 \\ \tilde{u}|_{\partial\Omega} & = \tilde{g}(x, t) \end{cases}$$

Bring the equation in the form

$$\begin{cases} -\nu \Delta u + \nabla p & = f(x, t) \\ \operatorname{div} u & = 0 \\ u|_{\partial\Omega} & = 0 \end{cases}$$

In order to solve this equation we need to introduce some new concepts

- Mixed Finite elements for variational formulation
- Inf-sup conditions

Mixed Finite elements, Variational Formulation

$$\begin{cases} -\nu \Delta u + \nabla p & = f(x, t) & | \cdot \phi \in X, \int_{\Omega} \\ \operatorname{div} u & = 0 & | \cdot \psi \in M, \int_{\Omega} \\ u|_{\partial\Omega} & = 0 \end{cases}$$

- Space for Velocity u : $X = H^1(\Omega)^2$
- Space for pressure $M = \{q \in L^2(\Omega) : \int_{\Omega} q dx = 0\}$

$$\begin{cases} a(u, \phi) + b(\phi, p) & = F(\phi) & \forall \phi \in X \\ b(u, \psi) & = 0 & \forall \psi \in M \end{cases}$$

with

$$a(v, \phi) = \nu \langle \nabla v, \nabla \phi \rangle, \quad b(\phi, \psi) = - \langle \operatorname{div} \phi, \psi \rangle$$

Inf-Sup, Existence and Uniqueness of the Solution

$$\begin{cases} a(u, \phi) + b(\phi, p) = F(\phi) & \forall \phi \in X \\ b(u, \psi) = 0 & \forall \psi \in M \end{cases}$$

Solve equations for $u \in Z = \{\phi \in X \mid b(\phi, \psi) = 0 \quad \forall \psi \in M\}$, that means solve

$$a(u, \phi) + \cancel{b(\phi, p)} = F(\phi) \quad \forall \phi \in Z$$

Solution exists due to coercivity and continuity of $a(u, \phi)$

Use u to solve

$$b(\phi, p) = F(\phi) - a(u, \phi) \quad \forall \phi \in X$$

Problem: ϕ and p are from different spaces:

Solution: use inf-sup-Condition $\sup_{v \in X} \frac{b(\phi, \psi)}{\|\phi\|_H} \geq \alpha \|\psi\|_M \quad \forall v \in X$
instead of coercivity

Relationship between Inf-Sup and Coercivity

Let $c(u, v)$ be an coercive bilinear form in a Hilbert space V , scalar product $\langle \cdot, \cdot \rangle_H$ and appropriate norm $\|\cdot\|_V$. then we have

$$\begin{aligned} c(u, v) &\geq \alpha \|u\|_H \|v\|_H \quad \forall u, v \in H \\ \implies \frac{c(u, v)}{\|v\|_H} &\geq \alpha \|u\|_H \quad \forall u, v \in H \\ \implies \sup_{v \in H} \frac{c(u, v)}{\|v\|_H} &\geq \alpha \|u\|_H \quad \forall u \end{aligned}$$

For existence and uniqueness of $b(\phi, p) = F(\phi) - a(u, \phi) \forall \phi \in X$

$$\implies \sup_{v \in X} \frac{b(\phi, \psi)}{\|\phi\|_H} \geq \alpha \|\psi\|_M \quad \forall v \in X$$

is enough.

Stokes Equations, Summary

Stokes equations	Poisson's equations
$\begin{cases} -\nu\Delta u + \nabla p & = f(x, t) \\ \operatorname{div} u & = 0 \\ u _{\partial\Omega} & = 0 \end{cases}$	$\begin{cases} -\nu\Delta u & = f(x, t) \\ u _{\partial\Omega} & = 0 \end{cases}$
Mixed FEM:	"not mixed" FEM
$\begin{cases} a(u, \phi) + b(\phi, p) & = F(\phi), \forall \phi \in X \\ b(u, \psi) & = 0, \forall \psi \in M \end{cases}$	$a(u, \phi) = F(\phi), \forall \phi \in X$
Inf-Sup	Coercivity
$\sup_{v \in X} \frac{b(\phi, \psi)}{\ \phi\ _H} \geq \alpha \ \psi\ _M, \forall v \in X$	$a(u, v) \geq \alpha \ u\ _X \ v\ _X$

Stokes Equations, Remarks to Numerics

$$\begin{cases} -\nu \Delta u + \nabla p & = f(x, t) \\ \operatorname{div} u & = 0 \\ u|_{\partial\Omega} & = 0 \end{cases}$$

- $\operatorname{div} u = 0$. X_h is not a subspace of the Solution!
 - Source of instabilities.
 - Construction of divergence preserving FEM-spaces of high order of precision is not simple
- Possible solutions:
 - Relaxation of the divergence free condition $\operatorname{div} u \approx \operatorname{div} u + \varepsilon u$
 - Special spaces like e.g. Crouzeix-Raviart for lower order precision

Stokes Equations simulations, Construction of Exact Solution

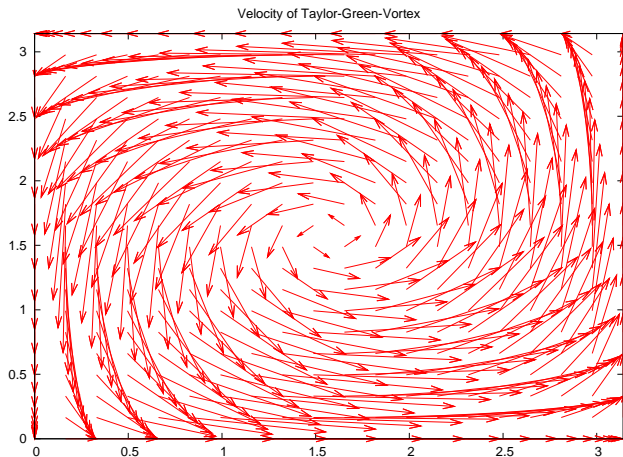
Taylor-Green Vortex

$$\begin{cases} \partial_t w + w \cdot \nabla w - \nu \Delta w + \nabla p & = 0 \\ \operatorname{div} w & = 0 \\ w|_{\partial\Omega} & = w|_{\partial\Omega} \end{cases}$$

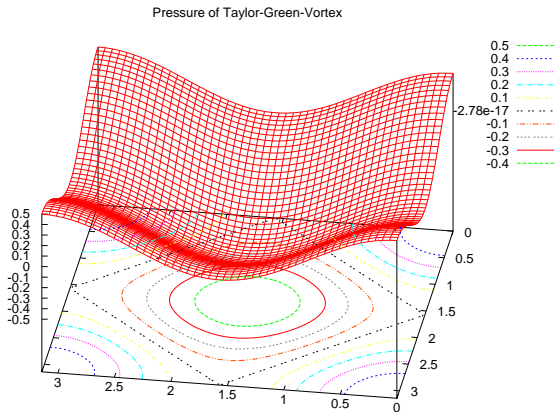
Stokes equations

$$\begin{cases} -\nu \Delta w + \nabla p & = \underbrace{-(\partial_t w + w \cdot \nabla w)}_{f(x)}|_{t=c} \\ \operatorname{div} w & = 0 \\ w|_{\partial\Omega} & = w|_{\partial\Omega, t=c} \end{cases}$$

Taylor-Green Vortex, Velocity



Taylor-Green Vortex, Pressure



Stokes Equations, Simulations

Mixed FEM: Taylor-Hood

$$X_h = \{\phi \in H^1([0, \pi])^2 : \phi|_T \in P^2 \forall T \in \mathcal{T}_h\}$$

$$M_h = \{\psi \in H^1([0, \pi]) : \psi|_T \in P^1 \forall T \in \mathcal{T}_h\}$$

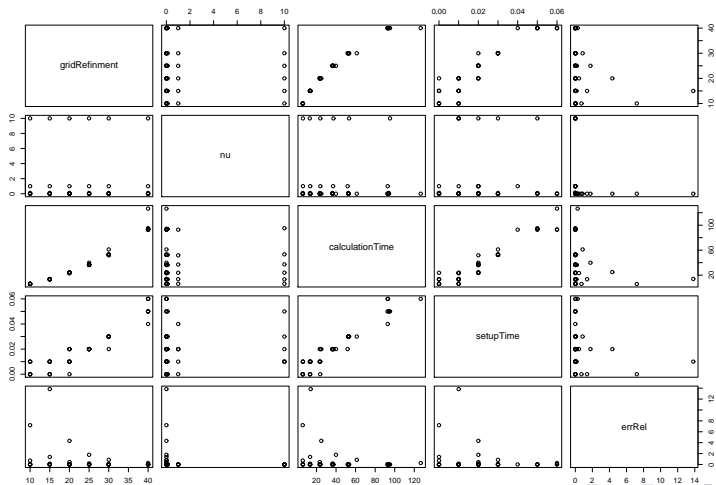
Mesh: $N \times N$ uniform mesh

Parameters:

ν	10	1.0	10^{-1}	10^{-2}	10^{-3}	10^{-4}	10^{-5}	10^{-6}	0
N	10	15	20	25	30	40			

Analytical reference: Taylor-Green Vortex

Stokes Equations, Postprocessing of Simulations

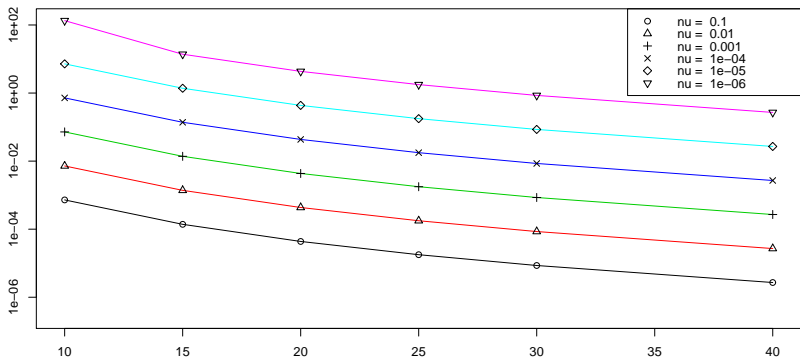


Stokes Equations, Postprocessing

Influence of grid refinement on relative error, Taylor-Hood

$$X_h = \{\phi \in H^1([0, \pi])^2 : \phi|_T \in P^2 \forall T \in T_h\},$$

$$M_h = \{\psi \in H^1([0, \pi]) : \psi|_T \in P^1 \forall T \in T_h\}$$



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Steady State Navier-Stokes Equations

$$\begin{cases} u \cdot \nabla u - \nu \Delta u + \nabla p & = f(x, t) \\ \operatorname{div} u & = 0 \\ u|_{\partial\Omega} & = g(x, t) \end{cases}$$

- Problem: Nonlinear part $u \cdot \nabla u$.

- Possible Solution: Iterative method

$$u^n \cdot \nabla u^{n+1} - \nu \Delta u^{n+1} + \nabla p^{n+1} = f(x, t) \quad (\text{Oseen Problem})$$

- Required: not too small ν , small f

- Possible reference, exact solution: Taylor-Green Vortex w

$$\begin{cases} w \cdot \nabla w - \nu \Delta w + \nabla p & = \underbrace{\partial_t w}_{f(x)}|_{t=c} \\ \operatorname{div} w & = 0 \\ w|_{\partial\Omega} & = w|_{\partial\Omega, t=c} \end{cases}$$

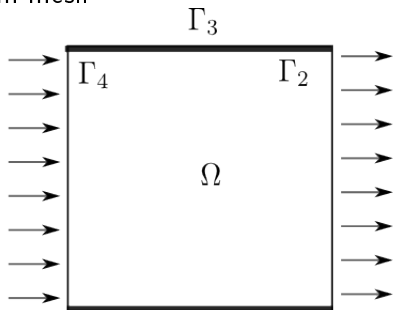
Steady State Navier-Stokes equations, Simulations

Mixed FEM: Taylor-Hood

$$X_h = \{\phi \in H^1([0, \pi])^2 : \phi|_T \in P^2 \forall T \in T_h\}$$

$$M_h = \{\psi \in H^1([0, \pi]) : \psi|_T \in P^1 \forall T \in T_h\}$$

Mesh: $N \times N$ uniform mesh



Boundary conditions $u_1|_{\Gamma_4} = c$ Γ_1 $u_1|_{\Gamma_2} = c$

Steady State Navier-Stokes equations.

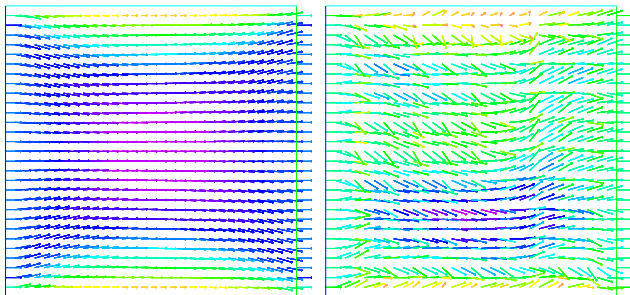


Figure:

Simulation with $k = 20$, grid= 10×10 : $\nu = 0.1$ left, $\nu = 10^{-7}$ right,
 $X_h = P2$, $M_h = P1$

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- Problem: Nonlinear part $u \cdot \nabla u$ and time derivative:
- Possible Solution: Discretization in time and linearisation (FreeFem++ *convex* operator)

$$\underbrace{\partial_t u + u^n \cdot \nabla u^n}_{\frac{1}{\Delta t}(u^{n+1} - u^n \circ X^n)} - \nu \Delta u^{n+1} + \nabla p^{n+1} = f^{n+1}$$

- Possible reference: Taylor-Green Vortex

$$\begin{cases} \partial_t w + w \cdot \nabla w - \nu \Delta w + \nabla p & = 0 \\ \operatorname{div} w = 0, & = 0 \end{cases}$$

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Navier Stokes equations with Relaxation

$$\begin{cases} \partial_t u + u \cdot \nabla u - \nu \Delta u + \nabla p & + \varepsilon \partial_{tt} u + \varepsilon \partial_t \nabla p = f(x, t) \\ \operatorname{div} u & = 0 \\ u|_{\partial\Omega} & = g(x, t) \quad \text{on } \partial\Omega \times [0, T] \end{cases}$$

- Problems: Time discretization $\varepsilon \partial_{tt} u + \varepsilon \partial_t \nabla p$
- Stabilities and Instabilities caused by ε parameters
- Possible solution: FD for time discretization special divergence preserving FEM spaces e.g. Crouzeix-Raviart
- Possible reference: Taylor-Green Vortex

$$\begin{cases} \partial_t w + w \cdot \nabla w - \nu \Delta w + \nabla p + \varepsilon \partial_{tt} w + \varepsilon \partial_t \nabla q & = \underbrace{\varepsilon \partial_{tt} w + \varepsilon \partial_t \nabla q}_{f(x,t)} \\ \operatorname{div} w = 0, & \end{cases}$$

Variational Formulation

$$\begin{aligned}
 0 \approx & \frac{1}{\Delta t} \langle u^{n+1}, \phi \rangle + \frac{\varepsilon}{\Delta t^2} \langle u^{n+1}, \phi \rangle + \nu \langle \nabla u^{n+1}, \nabla \phi \rangle \\
 & - \langle \operatorname{div} \phi, p^{n+1} \rangle + \frac{\varepsilon}{\Delta t} \langle \operatorname{div} \phi, p^{n+1} \rangle \\
 & - \frac{1}{\Delta t} \langle \operatorname{convect}(u^n, -\Delta t), \phi \rangle \\
 & + \frac{\varepsilon}{\Delta t^2} \langle -2u^n + u^{n-1}, \phi \rangle - \frac{\varepsilon}{\Delta t} \langle \operatorname{div} \phi, p^{n+1} \rangle \\
 & + \int_{\partial\Omega} \frac{\partial}{\partial n} u^{n+1} \phi ds + \int_{\partial\Omega} (\phi \cdot n) p ds \\
 & + \frac{\varepsilon}{\Delta t^2} \int_{\partial\Omega} \frac{\partial}{\partial n} (-2u^n + u^{n-1}) \phi ds - \frac{\varepsilon}{\Delta t} \int_{\partial\Omega} (\phi \cdot n) p^n ds
 \end{aligned}$$